

Characterization of Jet Induced Sidewall Pressures

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Final Report of SFW Cabin Noise Element

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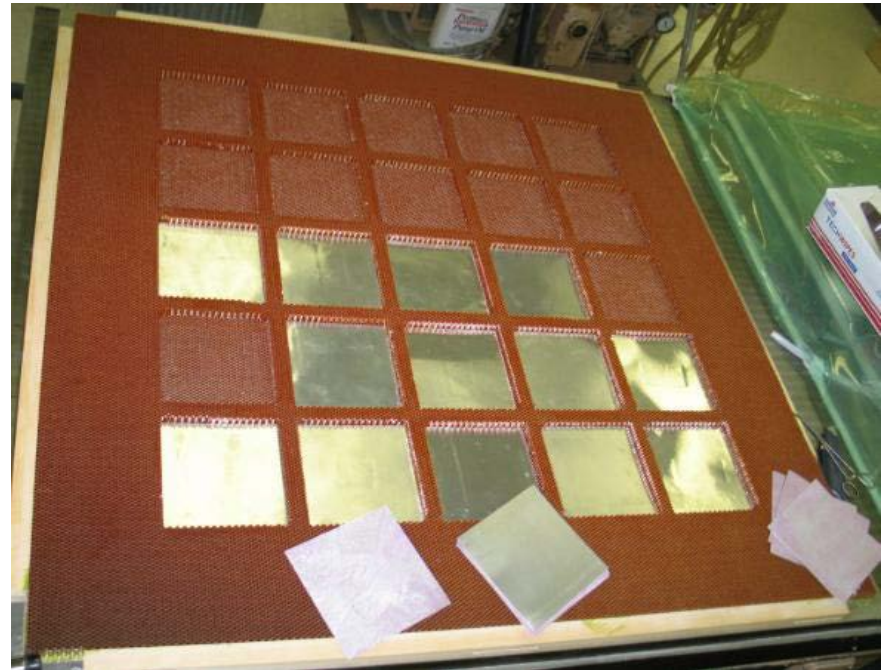
Outline

- Why Jet Noise is important to Cabin Noise
- Miller's extension of Ribner's two point correlation
- Palumbo/Miller extension of Harper-Bourne
- Results
- Concluding remarks

Path to Jet Noise Requirements

- Development of recessed/reinforced honeycomb composite
 - One patent awarded, another in process
 - INCE best paper award
 - Boeing interest

Optimization is necessary to fully realize potential of design



Path to Jet Noise Requirements

- Optimal design of structure requires accurate noise source models of sidewall **cross spectra**
 - Three sources were identified for characterization: TBL, Jet and Open Rotor
 - TBL completed in 2011
 - Jet noise milestone due in 2012
 - Open Rotor work will not be done

Summary of TBL Findings

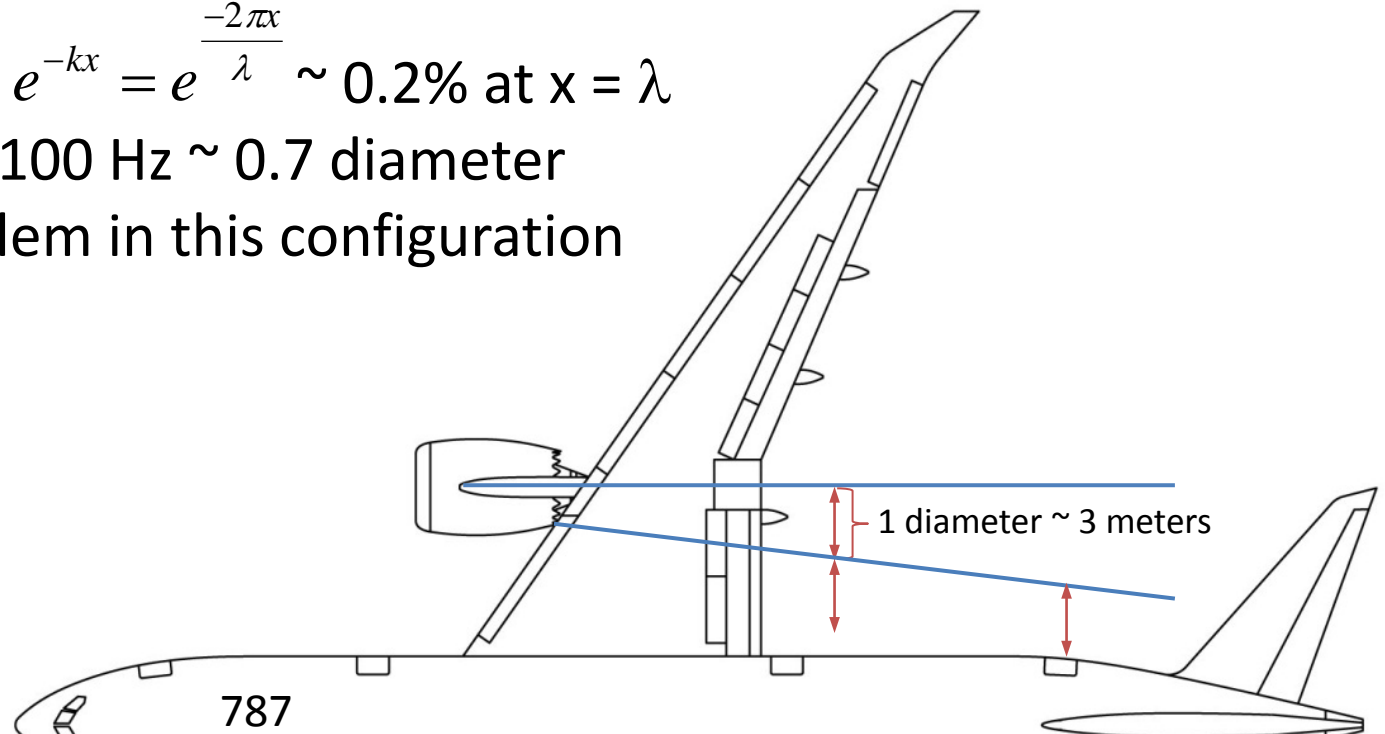
- Flight test raised more questions than it answered
- Estimate 7-8 dB uncertainty in noise radiation estimates based on measurements and theory
- 3 papers published in JSV



Aircraft Sidewall in Jet Near Field

- Negates some far field simplifications.
- Is evanescent field problematic?

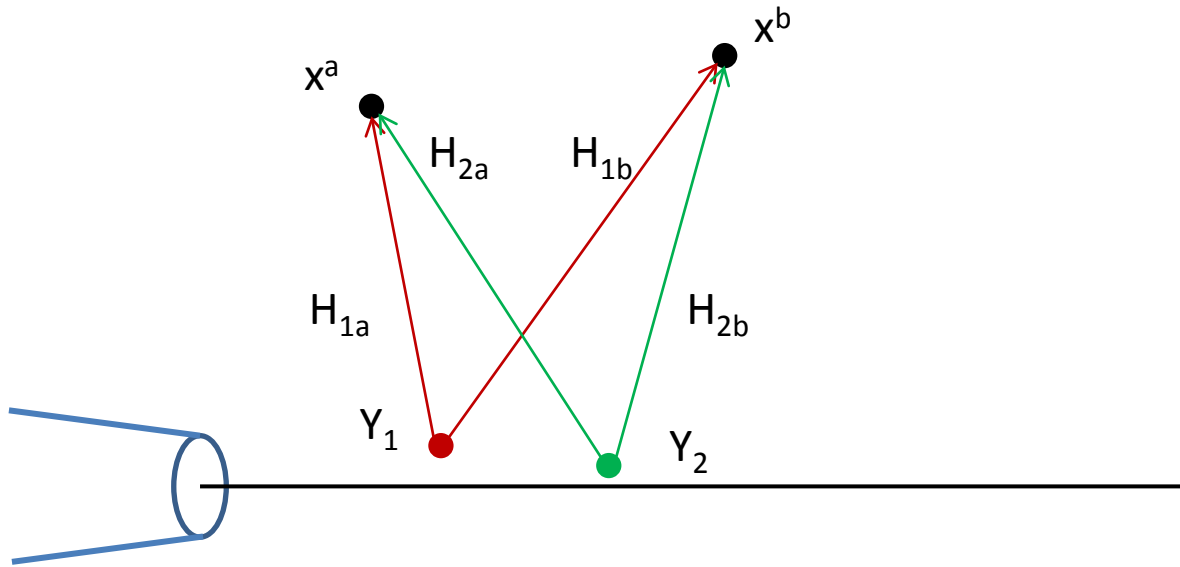
- decays like $e^{-kx} = e^{\frac{-2\pi x}{\lambda}} \sim 0.2\%$ at $x = \lambda$
- $\lambda \sim 2 \text{ m @ } 100 \text{ Hz} \sim 0.7 \text{ diameter}$
- Not a problem in this configuration



Relevant Jet Noise Models

- Nearly all jet noise models predict auto power spectrum in far field
- Ribner's model predicts far field correlation
 - Might be adaptable to near field
 - Has several limitations
- Harper-Bourne's model predicts near field cross power.
 - Requires an extensive experimental data set

Two Approaches



Miller/Ribner: Extend Ribner's two-point correlation theory using RANS solution and modern acoustic analogies.

Palumbo/Miller: Define source after Harper-Bourne and cross spectra in jet after Corcos and Leib. Propagate cross spectra to observers through analytically defined H.

Theory of Two-Point Correlations of Jet Noise, Herbert S. Ribner, 1976

- Far field approximation (radiation, geometry)
- No source convection, refraction or scattering
- Observer locations are equidistant
- Zero time delay between reception of signals
- Homogeneous, isotropic turbulence
- Exponential shape factor an oversimplification
- Source region non-compact and incoherent

Source Definition

Solution Simplification

Correlation Model

Ribner's Self Noise Prediction (Eq. 24)

$$R_{se}(\vec{x}^a, \vec{x}^b, \tau) = \underbrace{H(\vec{x}^a, \vec{x}^b)}_{\text{propagation}} \int_{jet} \underbrace{G_{yy}(\vec{Y})}_{\text{power}} \underbrace{e^{-2\omega_f |\tau + \Delta|}}_{\text{correlation}} d^3 \vec{Y}$$

where

R_{se} is correlation of acoustic pressure between equidistant points x^a and x^b at delay τ .

Y is the midpoint between two acoustic sources in the jet.

H is the transfer function of acoustic power between Y and x^a, x^b .

G_{yy} is the acoustic power at Y .

Δ is the time delay between the two acoustic sources at Y .

➤ Shear noise term has a similar form.

Extension of Prof. Ribner's Theory

$$\underbrace{G_{xy,\text{self}}(\vec{x}_1, \vec{x}_2, \omega)}_{\text{Cross-Spectra Self Noise}} = \int \underbrace{H(\vec{x}_1, \vec{x}_2, \vec{q})}_{\text{jet propagation}} \underbrace{G(\vec{x}_1, \vec{x}_2, \vec{q}, \omega)}_{\text{power}} \underbrace{C(\vec{x}_1, \vec{x}_2, \vec{q}, \omega)}_{\text{correlation}} d\vec{y}$$

where

q are the field-variables from a steady RANS solution

ω is the radial frequency

x_1 and x_2 are the observer locations

- Improvements: Modern source model, arbitrary observer placement, and freestream effects
- Require steady **RANS solution** and **experimental validation**
- Shear noise term has similar form
- Assumptions
 - Far-field
 - Statistical model of turbulent velocity covariance
 - Compact source
 - Isotropic turbulence

Prediction of Miller/Ribner to Experiment of Maestrello

Modified Ribner Model:

$$M_d = 1$$

$$M_j = 1$$

$$TTR = 3.20$$

$$D = 0.0508 \text{ m}$$

$$R/D = 100$$

Experiment by Maestrello:

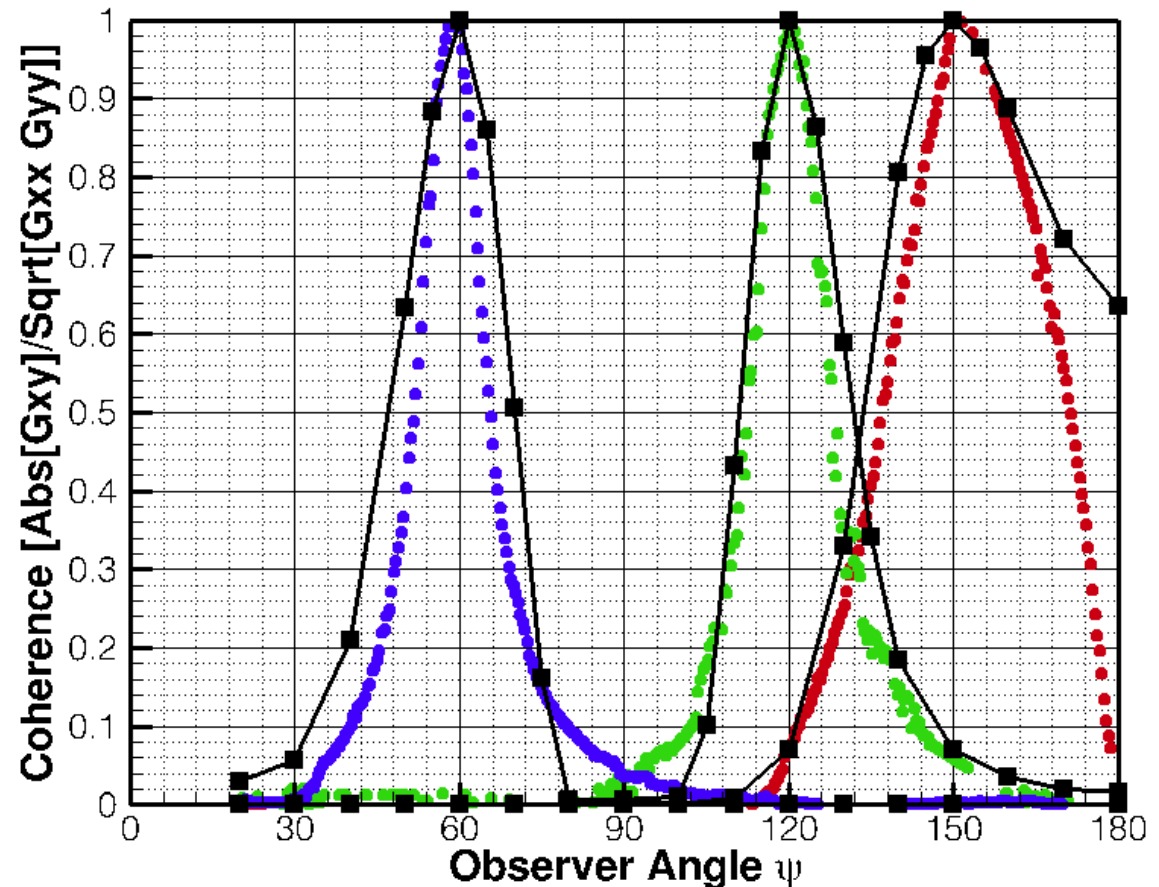
$$M_d = 1$$

$$M_j = 0.75$$

$$TTR = 1.00$$

$$D = 0.0254 \text{ m}$$

$$R/D = 169$$



Experiment: Maestrello, L., "Two-Point Correlations of Sound Pressure in the Far-Field of a Jet: Experiment," NASA TM X-72835, April 1976.

Palumbo, D. L. and Miller, S. A. E.

Harper-Bourne - Predicting the Jet Near-Field Noise of Combat Aircraft (2001)

- Semi-empirical
- Entropy noise assumed negligible
- Fixed frame of reference
- Self noise acoustically compact
- Observer located at $r >$ jet diameter
- Directivity and decay length taken as constant along jet axis

Harper-Bourne, M., "Predicting the Jet Near-Field Noise of Combat Aircraft," RTO AVT Symposium on Ageing Mechanisms and Control: Part A – Developments in Comp. Aero. & Hydro. Acous., Manchester, UK Oct. 8-11, 2001.

Harper-Bourne Cross Spectrum Prediction (Eq. 12)

$$G_{ab}(\vec{x}^a, \vec{x}^b, \omega) = \underbrace{G_{90}(\omega)}_{\text{power}} \int_y \underbrace{G_{yy}(y, \omega)}_{\text{power}} \underbrace{H(y, \vec{x}^a, \vec{x}^b, \omega)}_{\text{propagation \& correlation}} dy$$

$$H(y, \vec{x}^a, \vec{x}^b, \omega) = D(y, \vec{x}^a, \vec{x}^b, \omega) C_A(y, \vec{x}^a, \vec{x}^b, \omega) \frac{e^{-ik_a |\vec{x}^a - \vec{x}^b|}}{|\vec{x}^a| |\vec{x}^b|}$$

Where

G_{ab} is the cross spectrum of the pressures at x^a and x^b .

G_{90} is the power spectrum perpendicular to the jet axis

G_{yy} is the power in the jet at source location y

H is the transfer function that includes propagation and correlation

D is the cross power directivity

C_A is the cross power convective amplification

k_a is the wavenumber in air

➤ Correlation information is embedded in experimentally determined D and C_A .

Palumbo/Miller

Jet Cross Spectrum

$$G_{12}(y_1, y_2, \omega) = (G_r(y_1, \omega)G_r(y_2, \omega))^{1/2} e^{-|y_2 - y_1|/L} e^{-ik_j(y_2 - y_1)}$$

$$G_r(y, \omega) = G_{90}(\omega)G_{yy}(y, \omega)D_j^2$$

$$L = c_l D_j \left(1 - \text{Exp}\left[-c_s l_s St / D_j\right]\right) / St$$

Where

G_{12} is the cross spectrum in the jet

G_r is the power term relative to location y

L is the coherence length in the jet (after Morris and Boluriaan)

k_j is the wavenumber in the jet

D_j is the jet diameter

St is the Strouhal number

$$c_l = 7.16 \pi, c_s = 0.75, l_s = 0.2$$

➤ L and k_j should, most likely, be functions of jet axis position.

*From Morris and Boluriaan, "The Prediction of Jet Noise from CFD Data," AIAA Paper 2004-2977, 2004.

Palumbo/Miller Observer Cross Spectrum

$$G_{ab}(\vec{x}^a, \vec{x}^b, \omega) = \int \underbrace{H(y_1, \vec{x}^a, \omega)}_{\text{propagation}} \underbrace{G_{12}(y_1, y_2, \omega)}_{\text{power \& correlation}} \underbrace{H(y_2, \vec{x}^b, \omega)^*}_{\text{propagation}} d^2 y$$

$$H(y, \vec{x}, \omega) = D(\theta_o) C_A(\theta_o) \frac{e^{-ik_a r_o}}{4\pi r_o}$$

Where

θ_o is the angle from the jet axis to the observer

D is the directivity

C_A is the convective amplification

k_a is the wavenumber in air

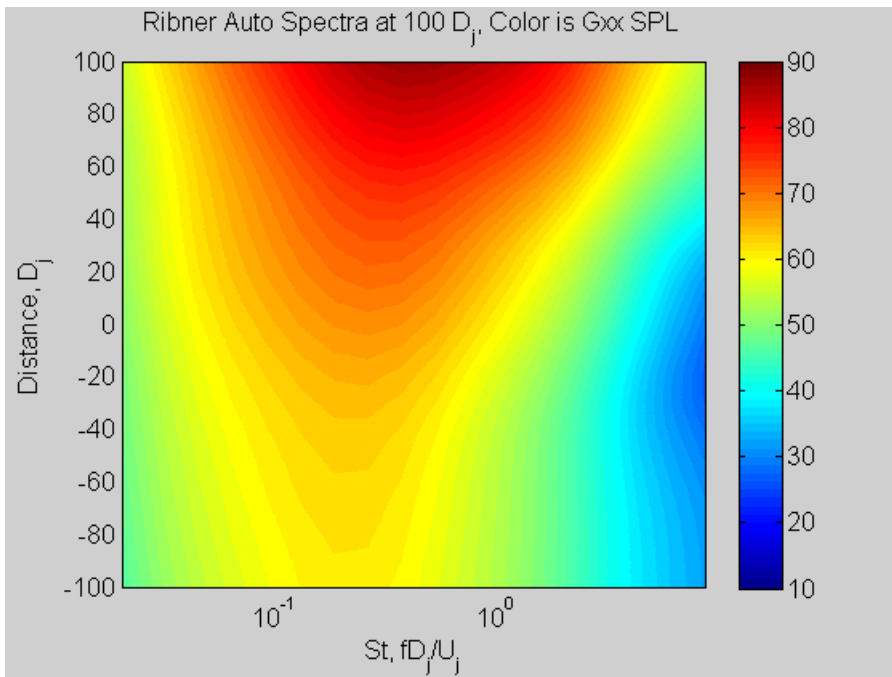
r_o is the distance from the source to the observer

Inherits limitations of Harper-Bourne model.

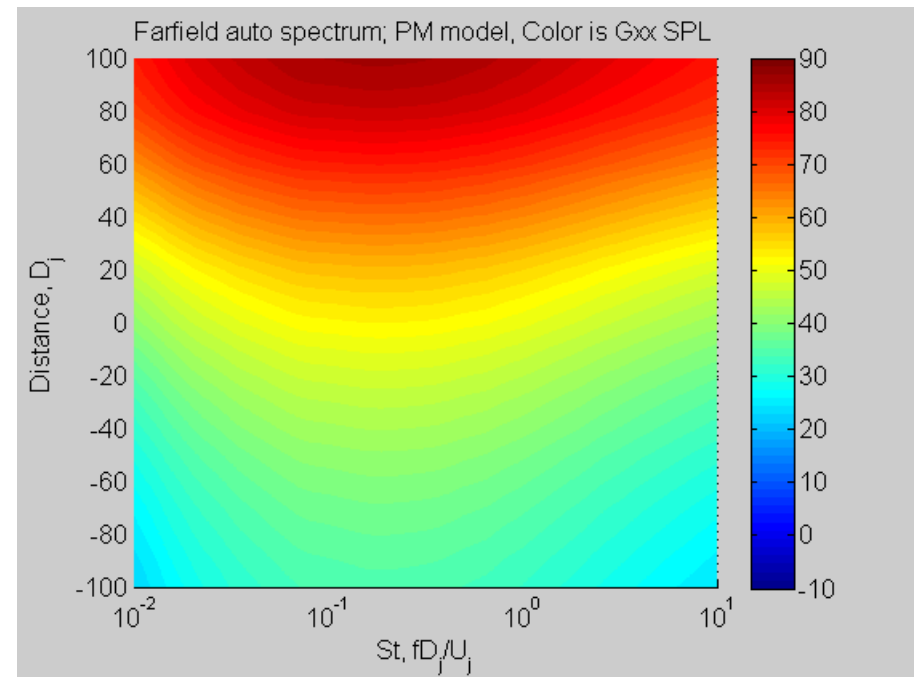
D and C_A should be function of jet axis position.

Far Field Auto Spectrum

Miller/Ribner, $y_2 = 100D$



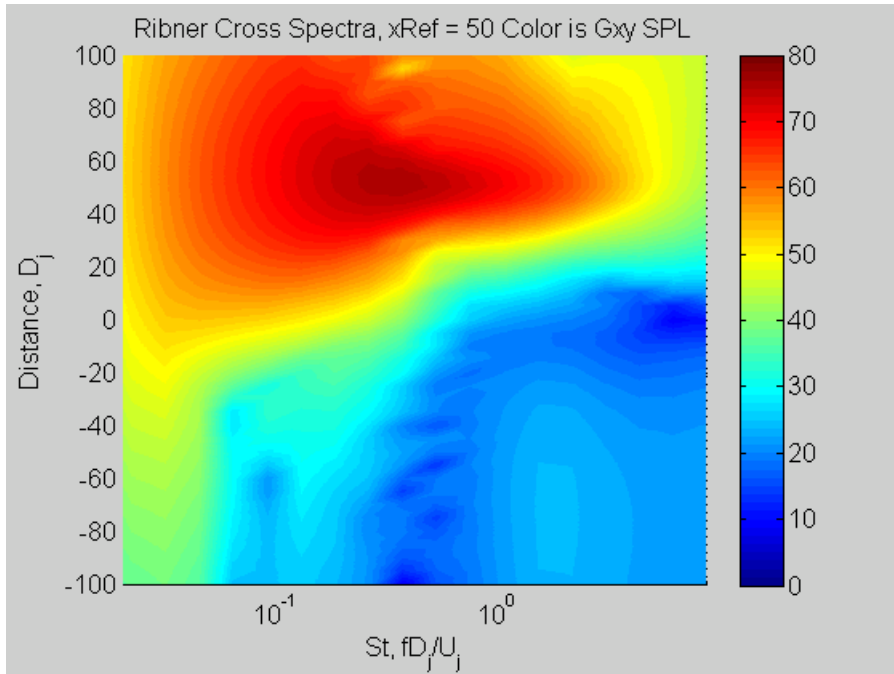
Palumbo/Miller, $y_2 = 100D$



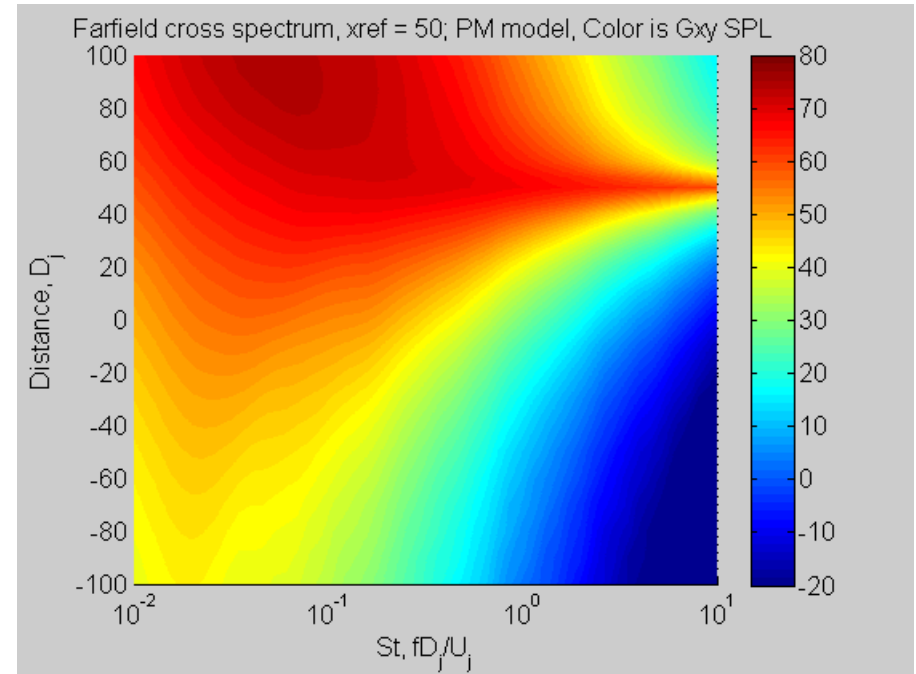
Palumbo/Miller (PM) auto spectrum levels are broader and peak at lower frequency.
The autopower in in PM model decays faster towards upstream direction

Far Field Cross Spectra

Miller/Ribner , $y_2 = 100D$



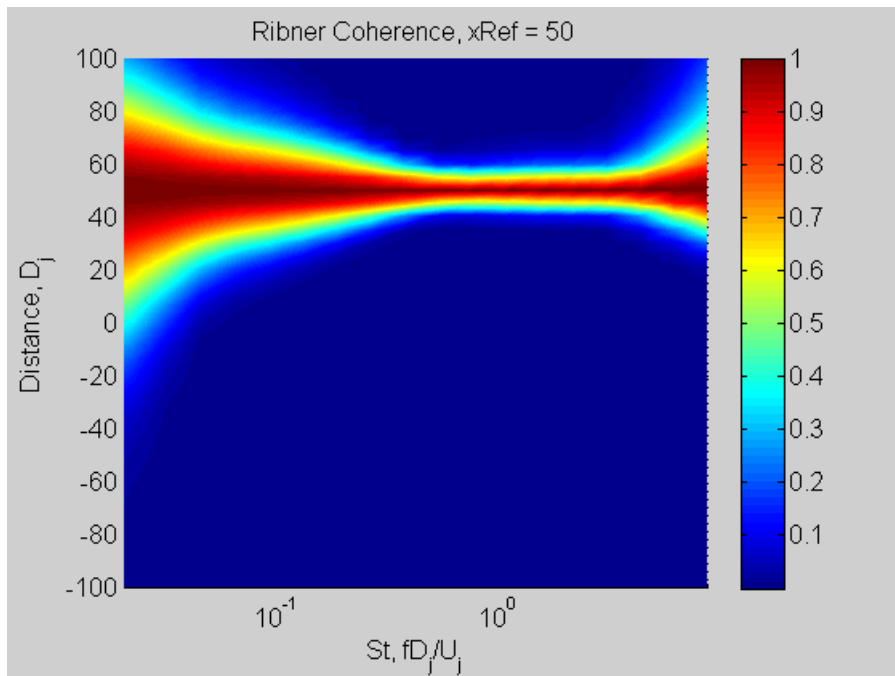
Palumbo/Miller , $y_2 = 100D$



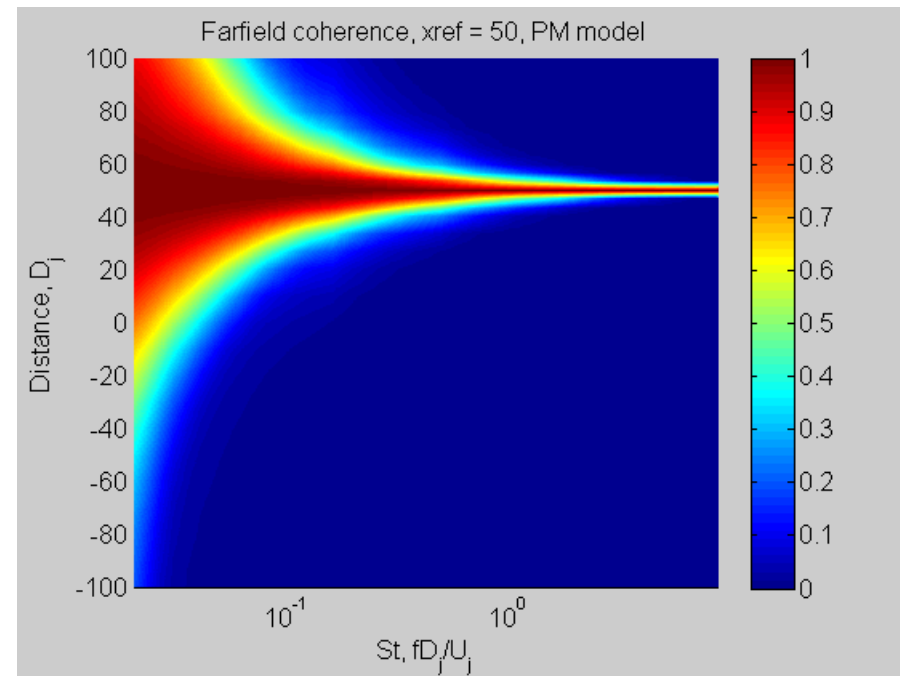
The PM model's peak in cross power occurs at a lower frequency and further downstream.

Far Field Coherence

Miller/Ribner , $y_2 = 100 D$



Palumbo/Miller , $y_2 = 100 D$

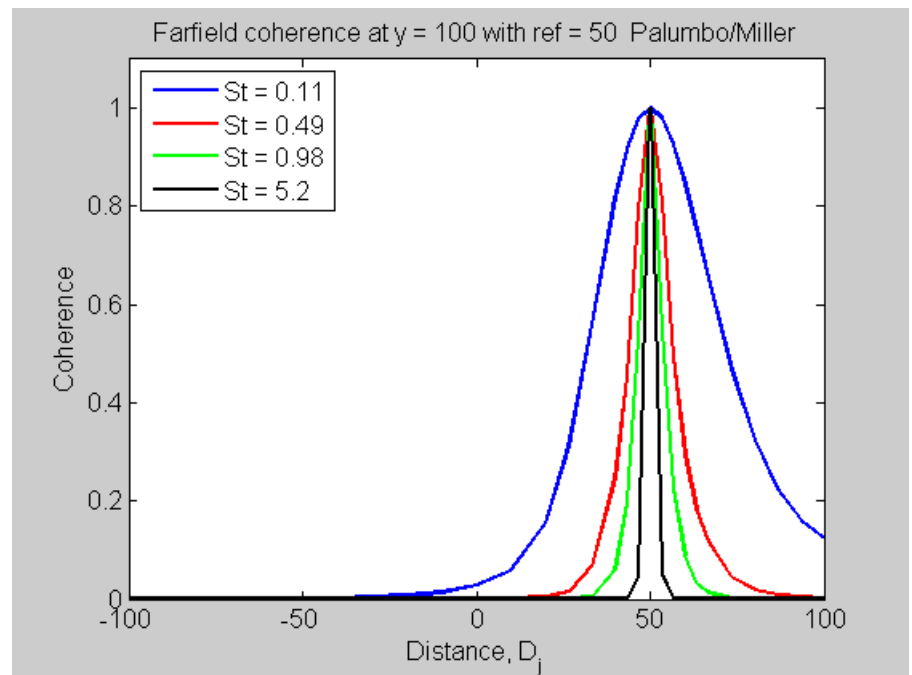
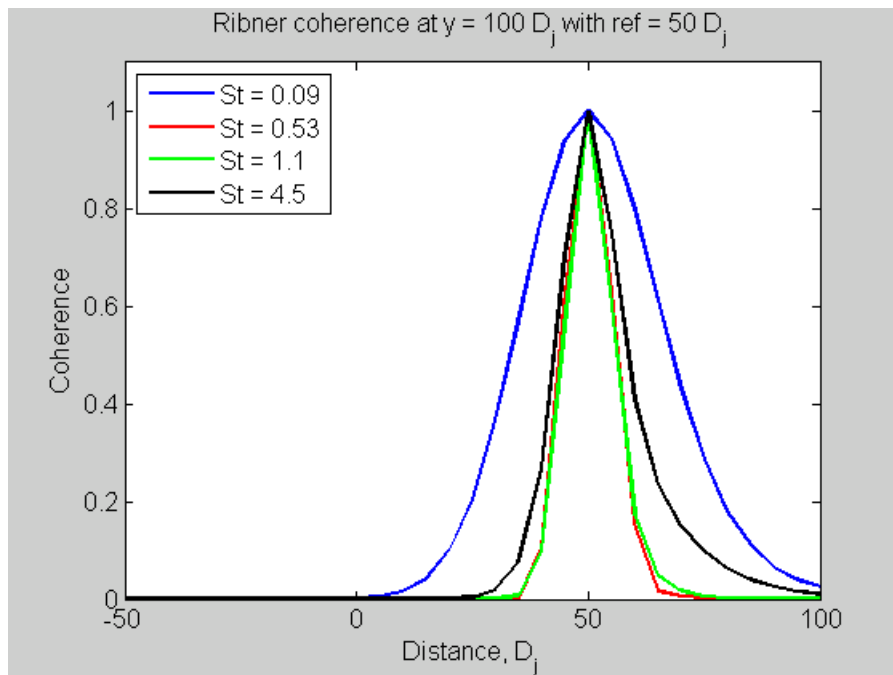


Shape of coherence is similar.

Far Field Coherence Cross Sections

Miller/Ribner , $y_2 = 100 D$

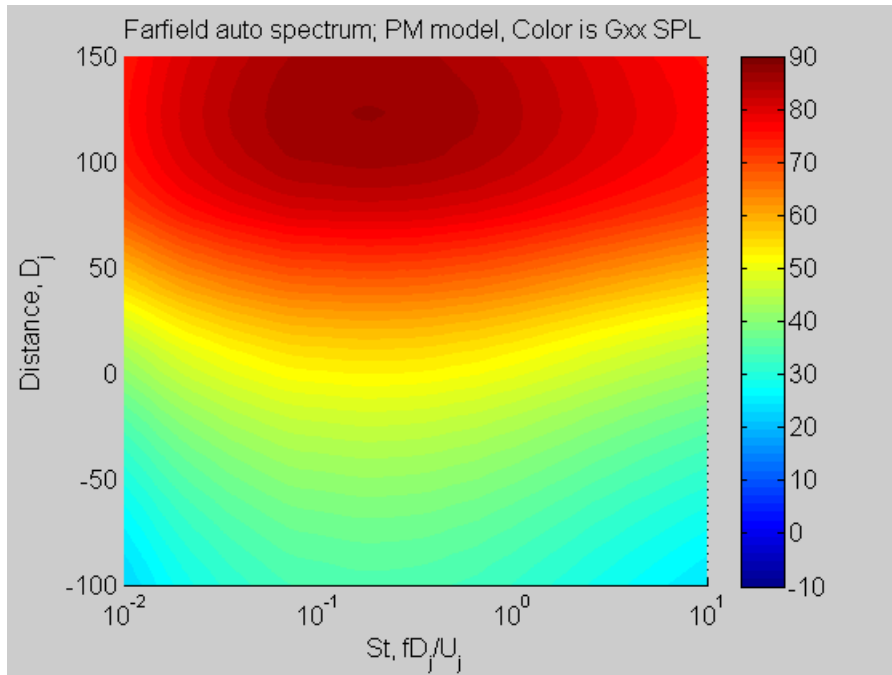
Palumbo/Miller , $y_2 = 100 D$



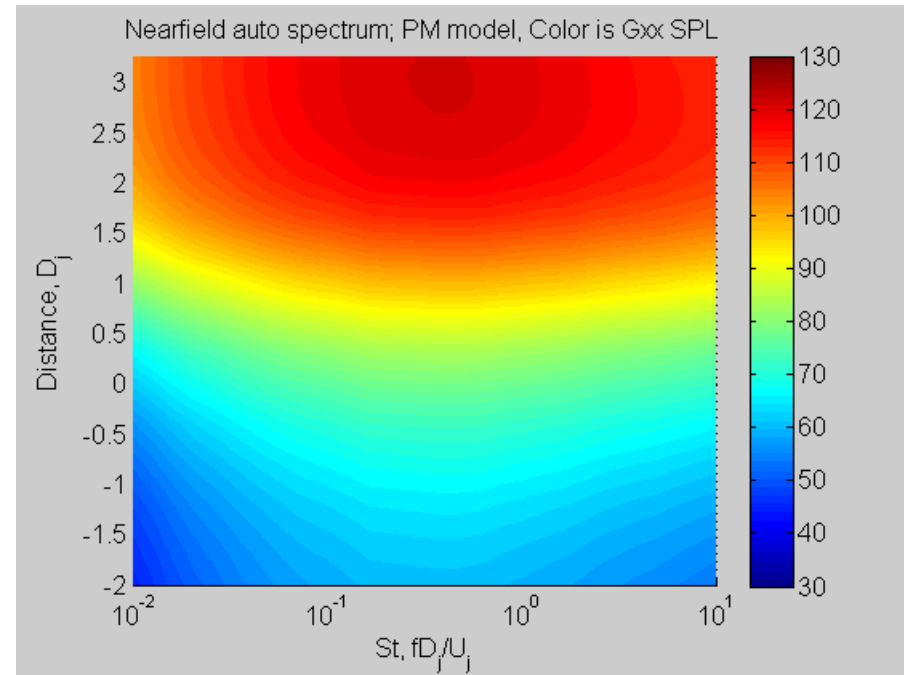
Shape of coherence similar at lowest frequency. Palumbo/Miller narrows quickly

Auto Spectrum: Far vs. Near Fields

Far Field, $y_2=100 D$



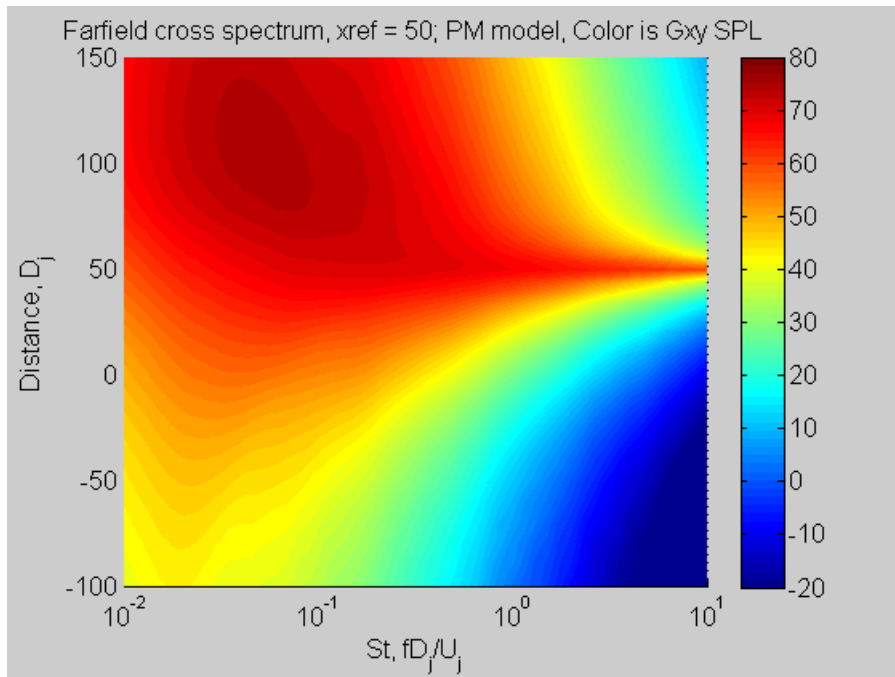
Near Field, $y_2=2 D$



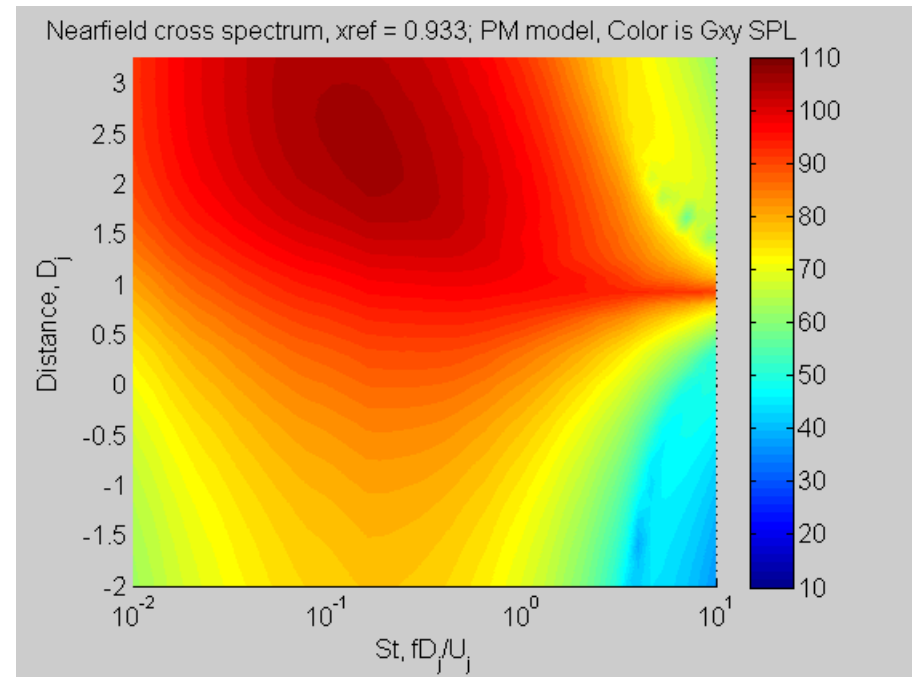
Directivity is relative to jet exit, so location of peak autopower is incorrect.

Cross Spectrum: Near vs. Far

Far Field, $y_2=100 D$



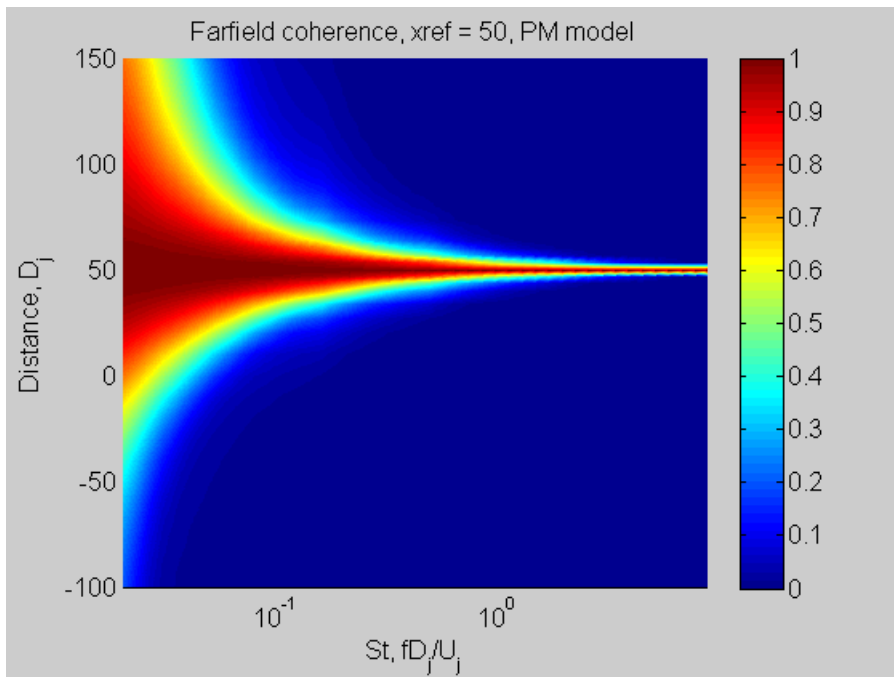
Near Field, $y_2=2 D$



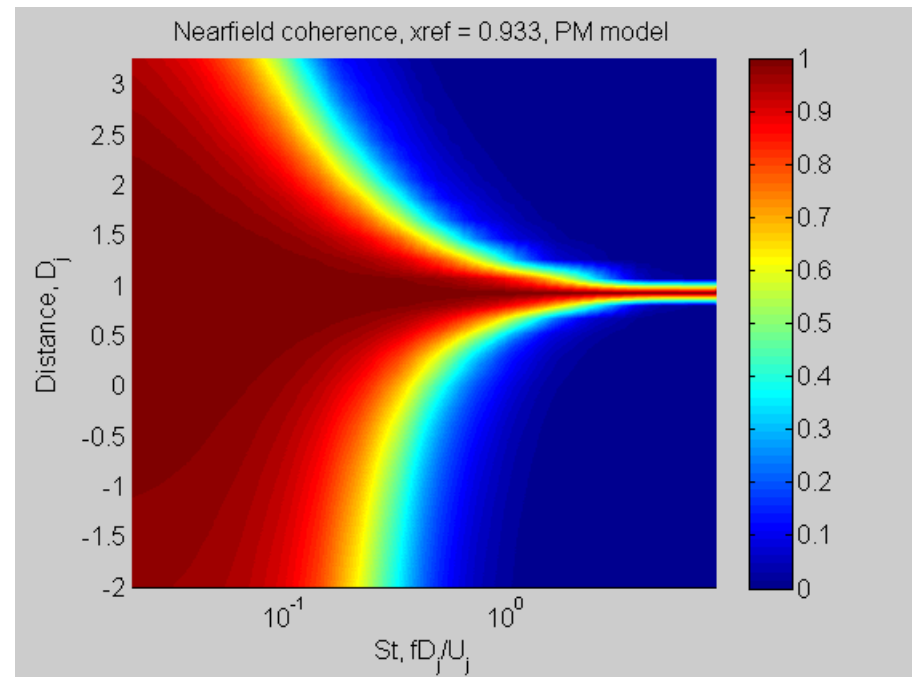
Peak power occurs at higher frequency, especially upstream.

Coherence: Near vs. Far

Far Field, $y_2=100 D$



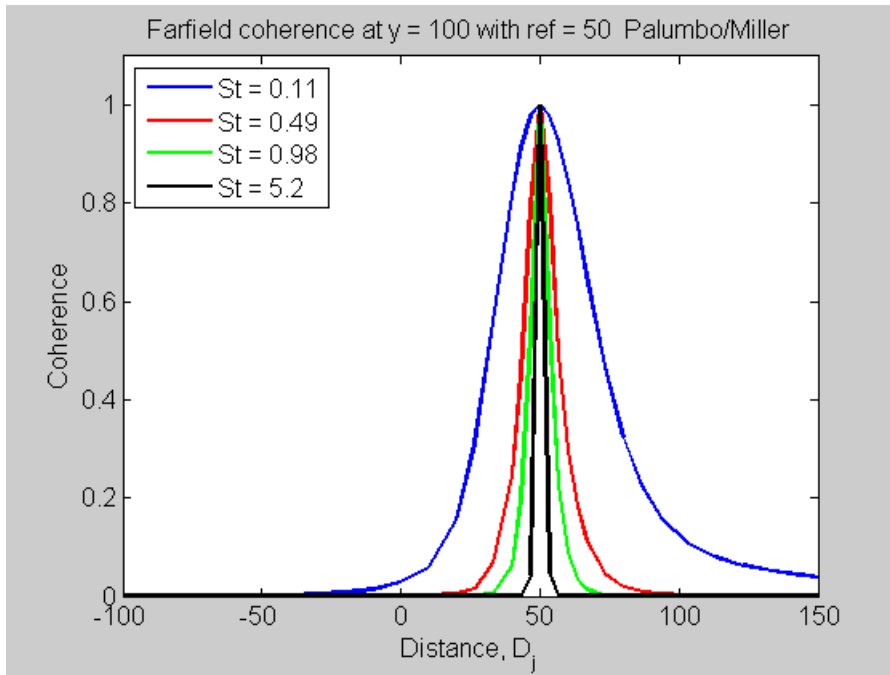
Near Field, $y_2=2 D$



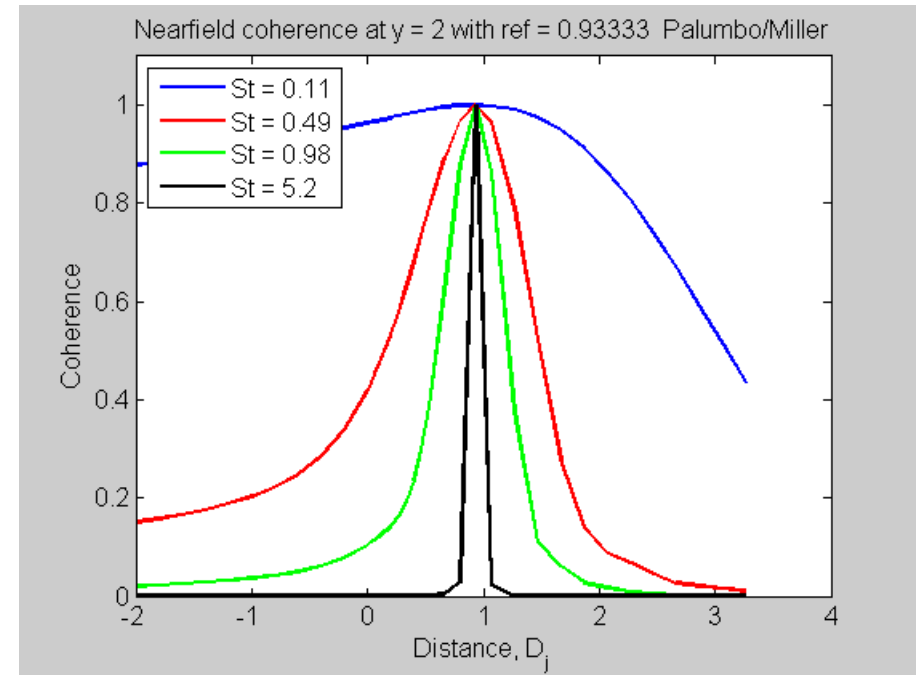
Coherence is much broader.

Coherence Cross Section: Near vs. Far

Far Field, $y_2=100 D$



Near Field, $y_2=2 D$



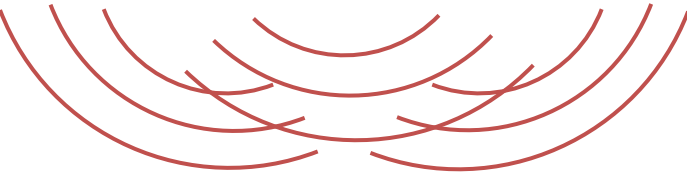
The very wide coherence pattern in the near field is made obvious in these cross sections

Re-radiation of Boundary Layer Excited Plate

Boundary layer at 150 m/s
and 300 m/s

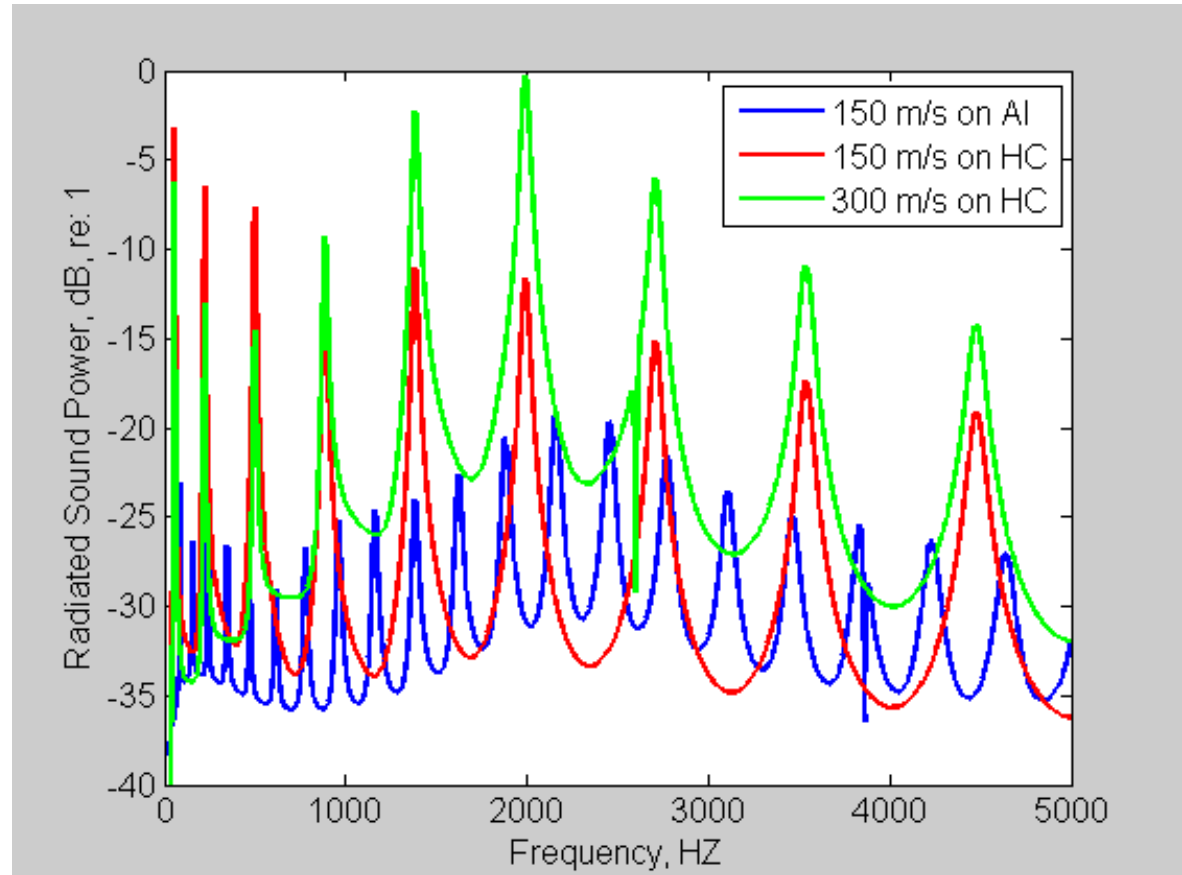


Aluminum and
Honeycomb Panels



Acoustic re-radiation

- Plate has both acoustic and hydrodynamic coincidence frequencies.
- When they occur in same frequency band, plate radiates very efficiently.



Conclusions

- Two approaches to formulating near-field predictions of the auto and cross spectra were pursued.
- The Miller/Ribner approach modernizes Ribner's theory but has yet to be extended to the near field.
- The Palumbo/Miller approach can be extended to the near field, but is relying on far field directivity and convective amplification terms.
- Both approaches require experimental data to fully mature the models.